



Galactic angular momenta and angular momentum correlations in the cosmological large-scale structure

Schäfer, 2009, Int. J. Mod. Phys. D18, 173

Stefano Camera

May 14, 2010



Outline

Structure formation

Tidal torquing

Angular momentum distribution

Angular momentum correlation

Implications for weak gravitational lensing

Structure formation

The formation of cosmic structures is based on gravitational amplification of seed perturbations in the cosmic density field.

Structure formation

The formation of cosmic structures is based on gravitational amplification of seed perturbations in the cosmic density field.

- 1 linearized continuity equation $a\dot{\delta} + \nabla \cdot \mathbf{v} = 0$
- 2 linearized Euler's equation $\dot{\mathbf{v}} + H\mathbf{v} = -\nabla\Phi/a$
- 3 Poisson's equation $\nabla^2\Phi = 4\pi G a^2 \rho_0 \delta$

Structure formation

The formation of cosmic structures is based on gravitational amplification of seed perturbations in the cosmic density field.

- ④ linearized continuity equation $a\dot{\delta} + \nabla \cdot \mathbf{v} = 0$
- ② linearized Euler's equation $\dot{\mathbf{v}} + H\mathbf{v} = -\nabla\Phi/a$
- ③ Poisson's equation $\nabla^2\Phi = 4\pi G a^2 \rho_0 \delta$

- $\delta(\mathbf{x}, a) = [\delta\rho(\mathbf{x}, a) - \rho_0(a)] / \rho_0(a)$ density contrast
- $H(a) = \dot{a}/a$ expansion history of the Universe
- $\Phi(\mathbf{x}, a)$ Newtonian potential

Structure formation

By decomposing the velocity field into the divergence $\theta = \nabla \cdot \mathbf{v}$ and the vorticity $\boldsymbol{\omega} = \nabla \times \mathbf{v}$

- ② linearized Euler's equation $\dot{\mathbf{v}} + H\mathbf{v} = -\nabla\Phi/a$

Structure formation

By decomposing the velocity field into the divergence $\theta = \nabla \cdot \mathbf{v}$ and the vorticity $\boldsymbol{\omega} = \nabla \times \mathbf{v}$

② linearized Euler's equation $\begin{cases} \dot{\theta} + H\theta + 4\pi G a \rho_0 \delta & = 0 \\ \dot{\boldsymbol{\omega}} + H\boldsymbol{\omega} & = 0 \end{cases}$

Structure formation

By decomposing the velocity field into the divergence $\theta = \nabla \cdot \mathbf{v}$ and the vorticity $\boldsymbol{\omega} = \nabla \times \mathbf{v}$

② linearized Euler's equation
$$\begin{cases} \dot{\theta} + H\theta + 4\pi G a \rho_0 \delta & = 0 \\ \dot{\boldsymbol{\omega}} + H\boldsymbol{\omega} & = 0 \end{cases}$$

$\boldsymbol{\omega}(a) \propto a^{-1}$, therefore any initial vortical excitation is wiped out rapidly.

Structure formation

By decomposing the velocity field into the divergence $\theta = \nabla \cdot \mathbf{v}$ and the vorticity $\boldsymbol{\omega} = \nabla \times \mathbf{v}$

② linearized Euler's equation
$$\begin{cases} \dot{\theta} + H\theta + 4\pi G a \rho_0 \delta & = 0 \\ \dot{\boldsymbol{\omega}} + H\boldsymbol{\omega} & = 0 \end{cases}$$

$\boldsymbol{\omega}(a) \propto a^{-1}$, therefore any initial vortical excitation is wiped out rapidly.

Galactic angular momenta cannot arise from vortical initial perturbations in the velocity field.

Tidal torquing

The current paradigm for generating the angular momentum of galaxies is tidal torquing, where the tidal gravitational field exerts a torquing momentum on the protogalactic object prior to collapse.

Tidal torquing

The current paradigm for generating the angular momentum of galaxies is tidal torquing, where the tidal gravitational field exerts a torquing momentum on the protogalactic object prior to collapse.

Lagrangian perturbation theory

$$\mathbf{x}(\mathbf{q}, t) = \mathbf{q} - D_+(t)\nabla\Psi(\mathbf{q})$$

$$\dot{\mathbf{x}}(\mathbf{q}, t) = -\dot{D}_+(t)\nabla\Psi(\mathbf{q})$$

- $D_+(t) = \delta(\mathbf{x}, a)/\delta(\mathbf{x}, 1)$ growth factor
- $\Psi(\mathbf{q})$ Zel'dovich displacement ($\nabla^2\Psi = \delta$)

Tidal torquing

The current paradigm for generating the angular momentum of galaxies is tidal torquing, where the tidal gravitational field exerts a torquing momentum on the protogalactic object prior to collapse.

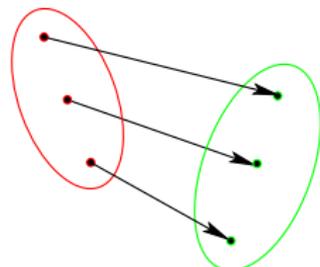
Lagrangian perturbation theory

$$\mathbf{x}(\mathbf{q}, t) = \mathbf{q} - D_+(t)\nabla\Psi(\mathbf{q})$$

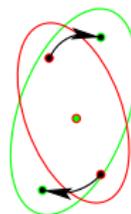
$$\dot{\mathbf{x}}(\mathbf{q}, t) = -\dot{D}_+(t)\nabla\Psi(\mathbf{q})$$

- $D_+(t) = \delta(\mathbf{x}, a)/\delta(\mathbf{x}, 1)$ growth factor
- $\Psi(\mathbf{q})$ Zel'dovich displacement ($\nabla^2\Psi = \delta$)

Euler frame



Lagrange frame



Euler frame

$$\mathbf{L} = \int d\mathbf{r} \rho(\mathbf{r})(\mathbf{r} - \bar{\mathbf{r}}) \times \mathbf{v} =$$

Euler frame

$$\begin{aligned}\mathbf{L} &= \int d\mathbf{r} \rho(\mathbf{r})(\mathbf{r} - \bar{\mathbf{r}}) \times \mathbf{v} = \\ &= \rho_0 a^5 \int d\mathbf{x} (\mathbf{x} - \bar{\mathbf{x}}) \times \dot{\mathbf{x}}\end{aligned}$$

- $\rho = \rho_0(1 + \delta) \simeq \rho_0$ because $\delta \ll 1$

Lagrange frame

Euler frame

$$\begin{aligned} \mathbf{L} &= \int d\mathbf{r} \rho(\mathbf{r})(\mathbf{r} - \bar{\mathbf{r}}) \times \mathbf{v} = \\ &= \rho_0 a^5 \int d\mathbf{x} (\mathbf{x} - \bar{\mathbf{x}}) \times \dot{\mathbf{x}} \end{aligned}$$

$$\begin{aligned} \mathbf{L} &= \rho_0 a^5 \int d\mathbf{q} (\mathbf{x} - \bar{\mathbf{x}}) \times \dot{\mathbf{x}} = \\ &\simeq \rho_0 a^5 \int d\mathbf{q} (\mathbf{q} - \bar{\mathbf{q}}) \times \dot{\mathbf{x}} = \end{aligned}$$

- $\rho = \rho_0(1 + \delta) \simeq \rho_0$ because $\delta \ll 1$

Lagrange frame

Euler frame

$$\begin{aligned} \mathbf{L} &= \int d\mathbf{r} \rho(\mathbf{r})(\mathbf{r} - \bar{\mathbf{r}}) \times \mathbf{v} = \\ &= \rho_0 a^5 \int d\mathbf{x} (\mathbf{x} - \bar{\mathbf{x}}) \times \dot{\mathbf{x}} \end{aligned}$$

$$\begin{aligned} \mathbf{L} &= \rho_0 a^5 \int d\mathbf{q} (\mathbf{x} - \bar{\mathbf{x}}) \times \dot{\mathbf{x}} = \\ &\simeq \rho_0 a^5 \int d\mathbf{q} (\mathbf{q} - \bar{\mathbf{q}}) \times \dot{\mathbf{x}} = \\ &= -\rho_0 a^5 \dot{D}_+ \int d\mathbf{q} (\mathbf{q} - \bar{\mathbf{q}}) \times (\mathbf{q} - \bar{\mathbf{q}}) \mathbf{H} \bar{\Psi} = \end{aligned}$$

- $\rho = \rho_0(1 + \delta) \simeq \rho_0$ because $\delta \ll 1$
- $\nabla\Psi \simeq \nabla\bar{\Psi} + (\mathbf{q} - \bar{\mathbf{q}})\mathbf{H}\bar{\Psi} \simeq (\mathbf{q} - \bar{\mathbf{q}})\mathbf{H}\bar{\Psi}$, with $\mathbf{H}\bar{\Psi}$ the tidal shear

Lagrange frame

Euler frame

$$\begin{aligned} \mathbf{L} &= \int d\mathbf{r} \rho(\mathbf{r})(\mathbf{r} - \bar{\mathbf{r}}) \times \mathbf{v} = \\ &= \rho_0 a^5 \int d\mathbf{x} (\mathbf{x} - \bar{\mathbf{x}}) \times \dot{\mathbf{x}} \end{aligned}$$

$$\begin{aligned} \mathbf{L} &= \rho_0 a^5 \int d\mathbf{q} (\mathbf{x} - \bar{\mathbf{x}}) \times \dot{\mathbf{x}} = \\ &\simeq \rho_0 a^5 \int d\mathbf{q} (\mathbf{q} - \bar{\mathbf{q}}) \times \dot{\mathbf{x}} = \\ &= -\rho_0 a^5 \dot{D}_+ \int d\mathbf{q} (\mathbf{q} - \bar{\mathbf{q}}) \times (\mathbf{q} - \bar{\mathbf{q}}) \mathbf{H} \bar{\Psi} = \\ &= \frac{1}{2} a^2 \dot{D}_+ [\mathbf{I}, \mathbf{H} \bar{\Psi}] \end{aligned}$$

- $\rho = \rho_0(1 + \delta) \simeq \rho_0$ because $\delta \ll 1$
- $\nabla \Psi \simeq \nabla \bar{\Psi} + (\mathbf{q} - \bar{\mathbf{q}}) \mathbf{H} \bar{\Psi} \simeq (\mathbf{q} - \bar{\mathbf{q}}) \mathbf{H} \bar{\Psi}$, with $\mathbf{H} \bar{\Psi}$ the tidal shear
- $\mathbf{I} = \rho_0 a^3 \int d\mathbf{q} (\mathbf{q} - \bar{\mathbf{q}})^T (\mathbf{q} - \bar{\mathbf{q}})$ inertia of the protogalactic cloud

$$\mathbf{X} \equiv \mathbf{I} \mathbf{H} \Psi$$

Angular momentum

$$\mathbf{L} = a^2 \dot{D}_+ \mathbf{X}^-$$

$$\mathbf{X} \equiv \mathbf{I} \mathbf{H} \Psi$$

Angular momentum

$$\mathbf{L} = a^2 \dot{D}_+ \mathbf{X}^-$$

Protogalactic objects will only acquire angular momentum if the eigensystems of their inertia and tidal shear are misaligned.

Angular momentum distribution

Galaxies form at local peaks in the large-scale structure by ellipsoidal collapse of the DM perturbation.

Angular momentum distribution

Galaxies form at local peaks in the large-scale structure by ellipsoidal collapse of the DM perturbation.

The peak criterion

- 1 $\delta(\mathbf{x}) > \nu$
- 2 $\nabla\delta(\mathbf{x}) = 0$
- 3 $\mathbf{H}\delta(\mathbf{x}) < 0$

Angular momentum distribution

Galaxies form at local peaks in the large-scale structure by ellipsoidal collapse of the DM perturbation.

The peak criterion

- ① $\delta(\mathbf{x}) > \nu$
- ② $\nabla\delta(\mathbf{x}) = 0$
- ③ $\mathbf{H}\delta(\mathbf{x}) < 0$

$$n_p(\nu) = \int d\mathbf{u} p(\mathbf{u}) \mathcal{C}(\mathbf{u})$$

- $p(\mathbf{u})d\mathbf{u}$ multivariate Gaussian random process
- $\mathcal{C}(\mathbf{u}) = \Theta[\delta - \nu] \delta_D[\nabla\delta] \Theta(-\lambda_i) |\lambda_1 \lambda_2 \lambda_3|$, with λ_i ($i = 1 \dots 3$) eigenvalues of $\mathbf{H}\delta$

n-th moment of the angular momentum

$$\langle \mathbf{L}^n \rangle = \frac{1}{n_p(\nu)} \int d\mathbf{u} p(\mathbf{u}) \mathcal{C}(\mathbf{u}) \mathbf{L}^n(\mathbf{u})$$

n-th moment of the angular momentum

$$\langle \mathbf{L}^n \rangle = \frac{1}{n_p(\nu)} \int d\mathbf{u} p(\mathbf{u}) \mathcal{C}(\mathbf{u}) \mathbf{L}^n(\mathbf{u})$$

The peak restriction makes the random process effectively non-Gaussian.

n -th moment of the angular momentum

$$\langle \mathbf{L}^n \rangle = \frac{1}{n_p(\nu)} \int d\mathbf{u} p(\mathbf{u}) \mathcal{C}(\mathbf{u}) \mathbf{L}^n(\mathbf{u})$$

The peak restriction makes the random process effectively non-Gaussian.

The moments of the angular momentum (e.g. the variance) decrease if peaks of higher overdensity are selected.

Angular momentum correlation

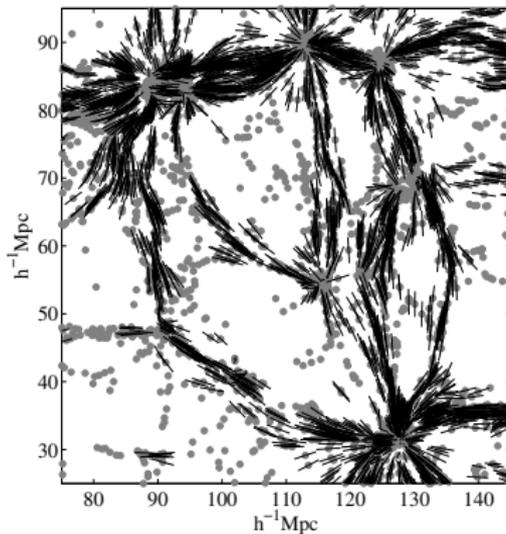
Linear régime

Correlation between neighbouring galaxies analytically via the covariance

$$\begin{aligned} C^{\mathbf{L}}(r) &\equiv \text{tr}[\langle \mathbf{L}(\mathbf{x})\mathbf{L}^T(\mathbf{x}') \rangle] = \\ &= a^4 \dot{D}_+^2 \text{tr}[\langle \mathbf{X}(\mathbf{x})\mathbf{X}(\mathbf{x}') \rangle - \langle \mathbf{X}(\mathbf{x})\mathbf{X}^T(\mathbf{x}') \rangle] \end{aligned}$$

Non-linear régime

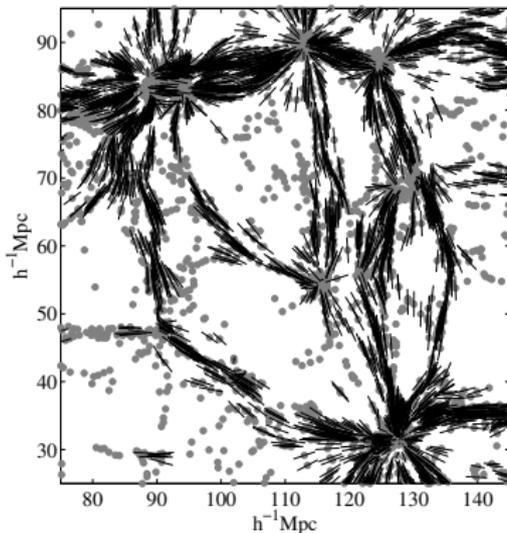
Correlation with the large-scale structure via N -body simulations



Hahn et al., 2007, MNRAS, 381, 41

Non-linear régime

Correlation with the large-scale structure via N -body simulations



Hahn et al., 2007, MNRAS, 381, 41

There is a correlation between galactic angular momenta (i.e. orientation of galaxies) and the large-scale structure.

Implications for weak gravitational lensing

	< 0	> 0
κ		
$\text{Re}[\gamma]$		
$\text{Im}[\gamma]$		

$$\mathbf{A} \equiv \begin{pmatrix} \kappa + \gamma_1 & \gamma_2 \\ \gamma_2 & \kappa - \gamma_1 \end{pmatrix} = \int d\chi W[\chi, n(\chi)] \nabla_{\perp} \Phi$$

Implications for weak gravitational lensing

	< 0	> 0
κ		
Re[γ]		
Im[γ]		

$$\mathbf{A} \equiv \begin{pmatrix} \kappa + \gamma_1 & \gamma_2 \\ \gamma_2 & \kappa - \gamma_1 \end{pmatrix} = \int d\chi W[\chi, n(\chi)] \nabla_{\perp} \Phi$$

Shear power spectrum

$$C^{\gamma\gamma}(\ell) = \frac{9}{16} H_0^4 \Omega_m^2 \ell^4 \int d\chi W^2[\chi, n(\chi)] \frac{D_+^2(a)}{a^2} P^{\delta}\left(\frac{\ell}{\chi}, \chi\right)$$

- $\gamma = \gamma_1 + i\gamma_2$ complex shear
- $\langle \delta_k(z) \delta_{k'}^*(z) \rangle = (2\pi)^3 \delta_D(\mathbf{k} - \mathbf{k}') P^{\delta}(k, z)$

Galaxy ellipticity

$$\epsilon \equiv \epsilon_+ + i\epsilon_\times$$

Uncorrelated intrinsic shapes are a common assumption in weak lensing

Galaxy ellipticity

$$\epsilon \equiv \epsilon_+ + i\epsilon_\times$$

Uncorrelated intrinsic shapes are a common assumption in weak lensing

- 1 Uncorrelated intrinsic ellipticities $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + \sigma_\epsilon^2/n$

Galaxy ellipticity

$$\begin{aligned}\epsilon &\equiv \epsilon_+ + i\epsilon_\times = \\ &= \frac{|\mathbf{L}|^2 - L_z^2}{|\mathbf{L}|^2 + L_z^2}\end{aligned}$$

Uncorrelated intrinsic shapes are a common assumption in weak lensing

- 1 Uncorrelated intrinsic ellipticities $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + \sigma_\epsilon^2/n$
- 2 Intrinsically correlated galaxy shapes $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + C^{\epsilon\epsilon}(\ell)$

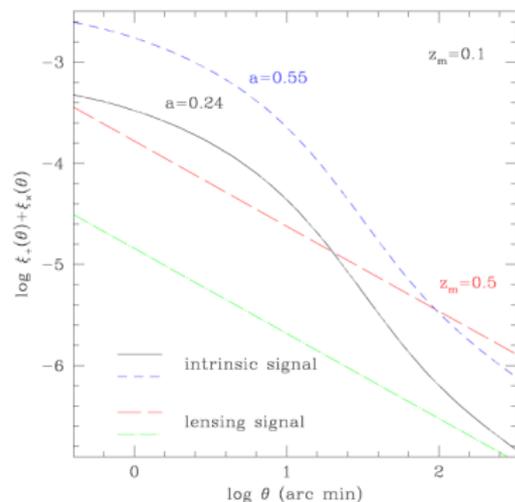
Galaxy ellipticity

$$\begin{aligned}
 \epsilon &\equiv \epsilon_+ + i\epsilon_\times = \\
 &= \frac{|\mathbf{L}|^2 - L_z^2}{|\mathbf{L}|^2 + L_z^2} = \\
 &\simeq \frac{\epsilon_{\text{int}} + \gamma}{1 + \gamma^* \epsilon_{\text{int}}}
 \end{aligned}$$

Uncorrelated intrinsic shapes are a common assumption in weak lensing

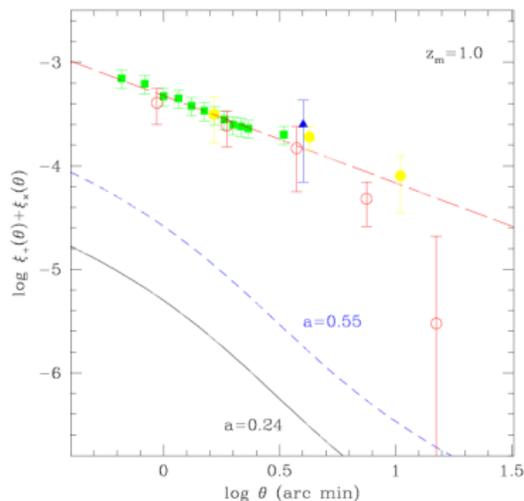
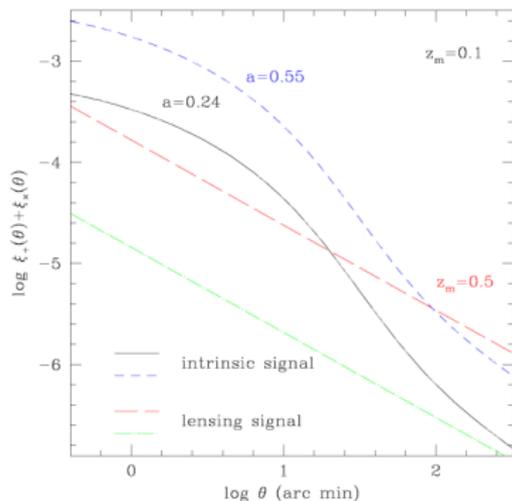
- ① Uncorrelated intrinsic ellipticities $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + \sigma_\epsilon^2/n$
- ② Intrinsically correlated galaxy shapes $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + C^{\epsilon\epsilon}(\ell)$
- ③ **Real case** $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + C^{\epsilon\epsilon}(\ell) + C^{\gamma\epsilon}(\ell)$

8 Real case $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + C^{\epsilon\epsilon}(\ell) + C^{\gamma\epsilon}(\ell)$



Crittenden et al., 2001, ApJ, 559, 552

8 Real case $\tilde{C}^{\gamma\gamma}(\ell) = C^{\gamma\gamma}(\ell) + C^{\epsilon\epsilon}(\ell) + C^{\gamma\epsilon}(\ell)$



Crittenden et al., 2001, ApJ, 559, 552

Thanks for your attention!